

# Gogny interactions with tensor terms

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First Gogny Conference

**Bruyères-le-Châtel, December 2015**

# Work in collaboration with

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- ▶ Shell evolution cannot be studied without tensor force (nuclei far from the stability line).
- ▶ Crucial in the study of properties of spin and spin-isospin states (Gamow-Teller and spin-dipole excitations).
- ▶ Some experiments which indicates the role of the tensor force:
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  2. Gaodefroy *et al.*, Phys. Rev. Lett **97**, 092501 (2006):  
Reduction of the neutron splitting  $1f_{7/2} - 2p_{3/2}$  in  $N = 28$  going from  $^{49}\text{Ca}$  to  $^{47}\text{Ar}$
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- ▶ Different fits have been done in order to fix the free parameters in each case:
  1. Adding a the tensor-isospin term, and modifying the strength of the spin-orbit term: D1ST
  2. Adding a pure tensor and tensor-isospin terms: D1ST2a, D1ST2b
  3. Adding a pure tensor, tensor-isospin and modifying the spin-orbit term: D1ST2c
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  1. Binding and single particle energies in HF approximation.
  2. Excitation states with DRPA and CRPA approximations:
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# Our Hartree-Fock (HF) approximation

- ▶ We consider as effective nucleon-nucleon interaction a finite-range two-body force of the type:

$$V(\vec{r}_1, \vec{r}_2) = \sum_{p=1}^6 V_p(\vec{r}_1, \vec{r}_2) O_p(1, 2) + V_{\text{SO}}(\vec{r}_1, \vec{r}_2) + V_{\text{DD}}(\vec{r}_1, \vec{r}_2) + V_{\text{Coul}}(\vec{r}_1, \vec{r}_2)$$

- ▶  $O_p(1, 2)$  indicates  $1, \vec{\tau}_1 \cdot \vec{\tau}_2, \vec{\sigma}_1 \cdot \vec{\sigma}_2, \vec{\sigma}_1 \cdot \vec{\sigma}_2 \vec{\tau}_1 \cdot \vec{\tau}_2, S_{12}, S_{12} \vec{\tau}_1 \cdot \vec{\tau}_2$ .
- ▶  $V_{\text{SO}}$  and  $V_{\text{DD}}$ , terms of zero-range (like the corresponding terms in Skyrme-like forces)
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- ▶ We solve, in coordinate space, a set of equations of the type:

$$-\frac{\hbar^2}{2m_k} \nabla_1^2 \phi_k(\vec{r}_1) + U(\vec{r}_1) \phi_k(\vec{r}_1) - \int d^3 r_2 W(\vec{r}_1, \vec{r}_2) \phi_k(\vec{r}_2) = \epsilon_k \phi_k(\vec{r}_1)$$

- ▶ Hartree (Direct) term
- ▶ Fock (Exchange) term

## Fit of the tensor terms: D1ST interaction

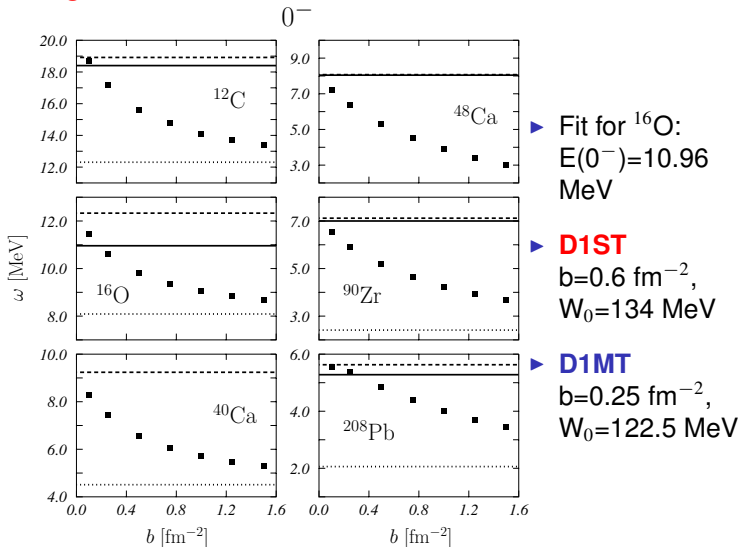
$$v_{6,b}(r) = v_{6,AV18}(r) \left(1 - e^{-br^2}\right)$$

$$V_6(q)S_{12}(\mathbf{q}) = \int d^3r e^{i\mathbf{q}\cdot\mathbf{r}} v_6(r) S_{12}(\mathbf{r}) = -4\pi \int dr r^2 j_2(qr) v_6(r) S_{12}(\mathbf{r})$$



# Fit of the tensor force: D1ST

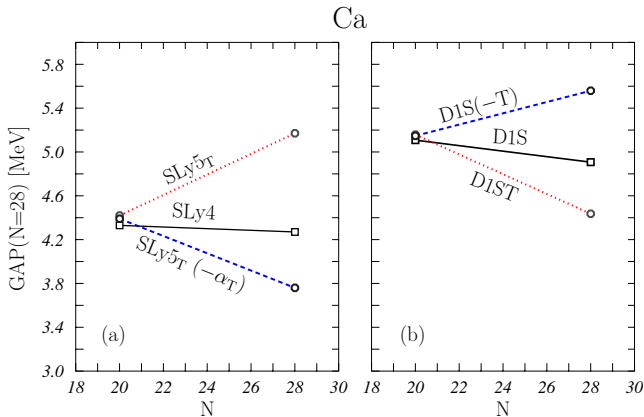
## Energies of the first $0^-$ states



Phenomenological RPA with LM +  $v_{6,b}(r)$

# Fit of the tensor force: D1ST2a and D1ST2b

- Experimentally, the difference between the energies of the single-particle neutron  $2p_{3/2}$ ,  $1f_{7/2}$  states increases from  $^{40}\text{Ca}$  to  $^{48}\text{Ca}$ : O. Sorlin and M.-G. Pourquet, Prog. Part. Nucl. Phys. **61**, 602 (2008)



# Fit of the tensor force: D1ST2a and D1ST2b

N. Onishi and J.W. Negele, NPA301 (1978), 336

$$\begin{aligned} V_{\text{tensor}}(\vec{r}_1, \vec{r}_2) &= (V_{T1} + V_{T2} P_{12}^T) S_{12} e^{-(r_1-r_2)^2/\mu_T^2} \\ &= \left[ \left( V_{T1} + \frac{1}{2} V_{T2} \right) + \frac{1}{2} V_{T2} \boldsymbol{\tau}(1) \cdot \boldsymbol{\tau}(2) \right] S_{12} e^{-(r_1-r_2)^2/\mu_T^2} \end{aligned}$$

**D1ST2a** → neutron  $1f$  splitting in  $^{48}\text{Ca}$  and  $0^-$  state of  $^{16}\text{O}$ :

$$V_{T1} = -135 \text{ MeV}, V_{T2} = 115 \text{ MeV}$$

**D1ST2b** →  $N = 28$  neutron gap increase from  $^{40}\text{Ca}$  to  $^{48}\text{Ca}$  as obtained in HF calculations with the SLy5T force and  $0^-$  state of  $^{16}\text{O}$ :

$$V_{T1} = -182 \text{ MeV}, V_{T2} = 102 \text{ MeV}$$

# Fits of the tensor force: D1ST2c

Following the strategy of Zalewski *et al.*, Phys. Rev. **C77**, 024316 (2008):

**D1ST2c** → neutron  $1f$  splitting in  $^{40}\text{Ca}$ ,  $^{48}\text{Ca}$  and  $^{56}\text{Ni}$

1. First, we fit the splitting  $1f$  in  $^{40}\text{Ca}$  by modifying the spin-orbit parameter  $W_{\text{LS}}$ ,
2. second, we fit the splitting  $1f$  in  $^{48}\text{Ca}$  adjusting the like-particle part of the Gogny tensor term and,  $V_{T1} + V_{T2}$ ,
3. finally, we use the  $^{56}\text{Ni}$  to fit the neutron-proton contribution of the tensor term,  $V_{T2}$ .<sup>1</sup>

$$W_{\text{LS}} = 103 \text{ MeV fm}^5, \quad V_{T1} = -135 \text{ MeV}, \quad V_{T2} = 60 \text{ MeV}$$

**D1MT2c** → following the same procedure, and using D1M as starting point we have fit another interaction

$$W_{\text{LS}} = 95 \text{ MeV fm}^5, \quad V_{T1} = -175 \text{ MeV}, \quad V_{T2} = 40 \text{ MeV}$$

---

<sup>1</sup>M. Grasso and M. A, Phys. Rev. **C88**, 054328 (2014)

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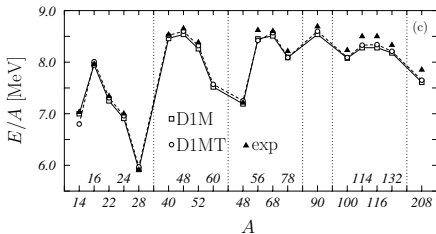
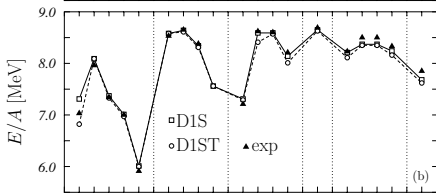
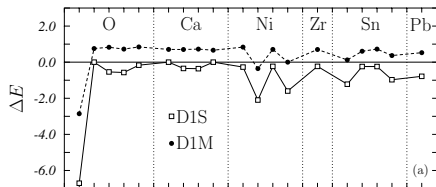
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# Bulk properties of spherical nuclei

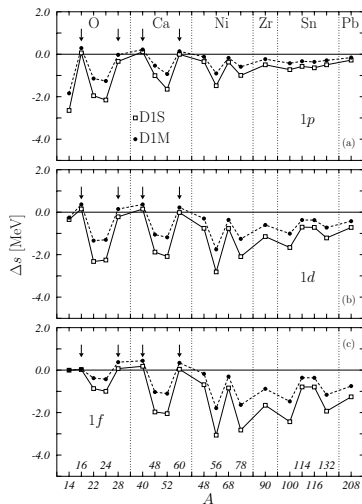


$$\Delta E = 100 \frac{E_{D1\alpha T} - E_{D1\alpha}}{E_{D1\alpha}}$$

$\alpha \equiv S, M$

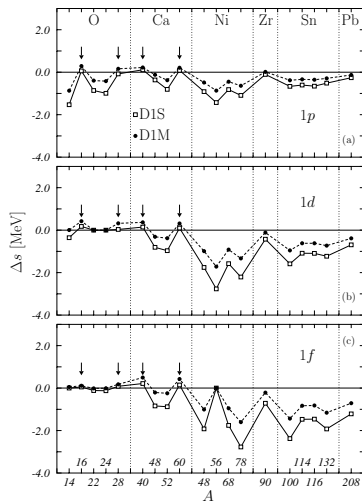
# s.p. energies of spin-orbit partners

**Protons**



$$S = \epsilon_{l-1/2} - \epsilon_{l+1/2}$$

**Neutrons**



$$\Delta s = s_{D1\alpha T} - s_{D1\alpha}$$



# Discrete RPA: $0^-$ states

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	exp	D1S	D1ST	D1M	D1MT
$^{12}\text{C}$	18.40	19.63	14.42	18.83	15.27
$^{16}\text{O}$	10.96	13.95	10.94	13.08	10.96
$^{40}\text{Ca}$	10.78	12.22	9.57	11.56	9.60
$^{48}\text{Ca}$	8.05	14.10	11.63	12.85	11.26
$^{208}\text{Pb}$	5.28	8.27	7.93	8.24	7.92

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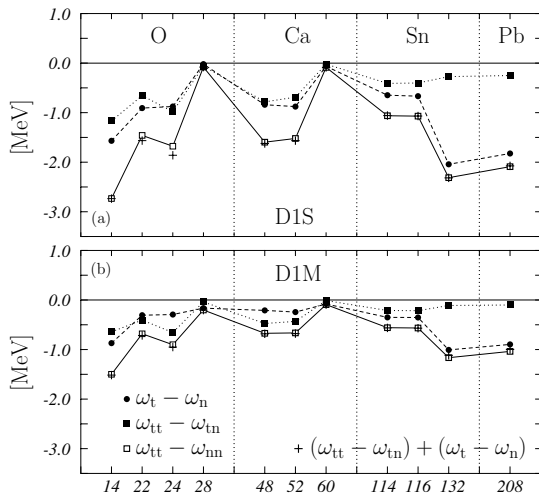
M. A *et al.* Phys. Rev. **C83** (2011) 064306

Exp. values:

A. Heusler *et al.* Phys. Rev. **C75** (2007) 024312

<http://www.nndc.bnl.gov/>

# Discrete RPA: $1^+$ excitations in $N \neq Z$ nuclei

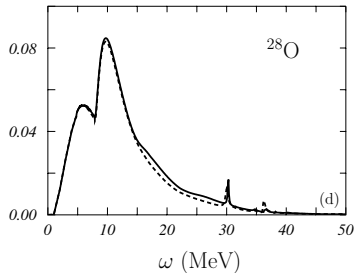
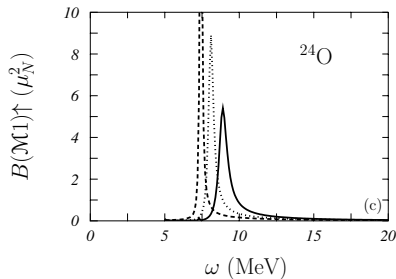
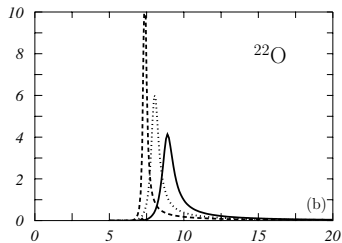
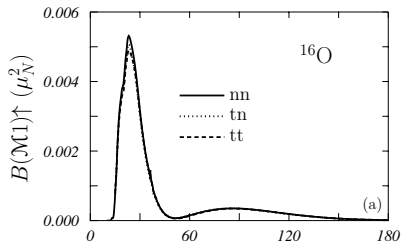


First  $1^+$  state in  $^{208}\text{Pb}$

	$E(1^+)$ [MeV]	$B(M1)_1$ [ $\mu_n^2$ ]
exp	5.85	2.0
D1S	7.80	5.08
D1ST	4.76	2.41
D1M	6.50	2.33
D1MT	4.82	1.80

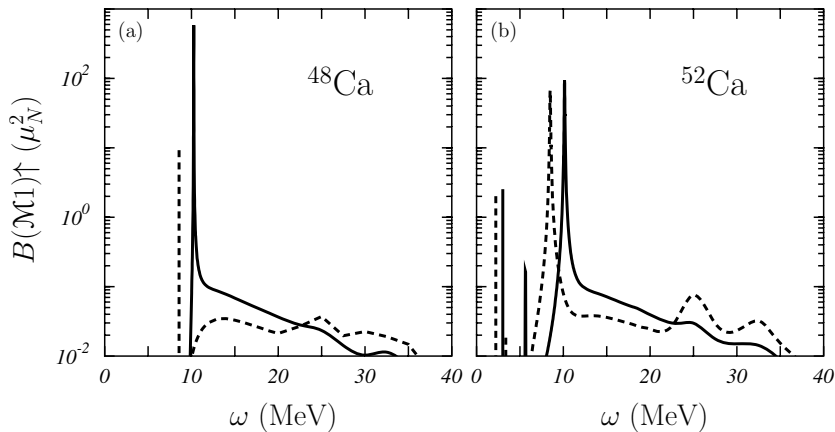
M. A *et al.* Phys. Rev. **C83** (2011) 064306

# Continuum RPA: Magnetic dipole response in O isotopes



main p-h excitation  $\Rightarrow$   $[(\nu 1d_{3/2})(\nu 1d_{5/2})^{-1}]$

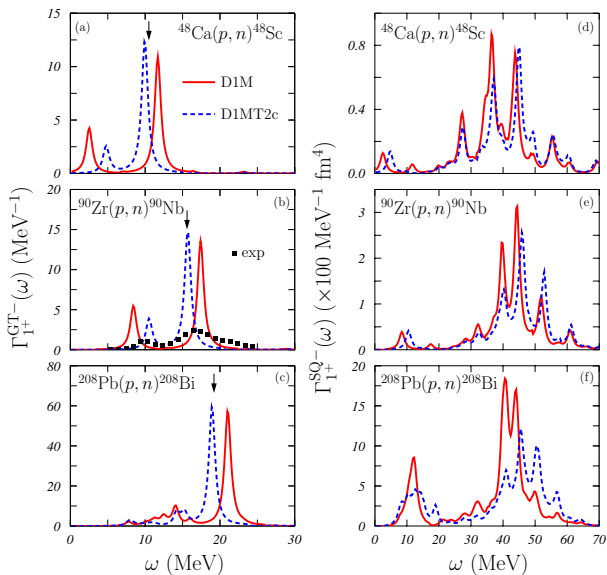
# Continuum RPA: Magnetic dipole response in Ca isotopes



main p-h excitation  $\Rightarrow [(\nu 1f_{5/2})(\nu 1f_{7/2})^{-1}]$

$E(1^+)$ : 10.23 MeV (EXP)    10.15 MeV (D1S)    8.56 MeV (D1ST)

# Charge exchange excitations: $1^+$ GT and SQ



# Splitting in $^{40}\text{Ca}$ , $^{36}\text{S}$ and $^{34}\text{Si}$ : $N = 20$ isotones

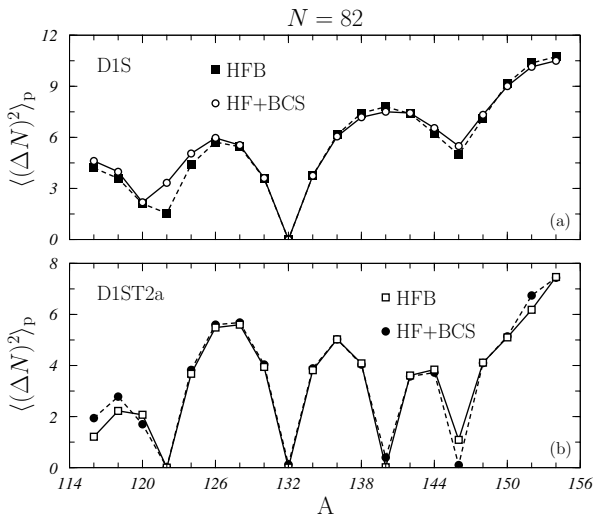
## Tensor induced and pure spin-orbit effects

	From $^{40}\text{Ca}$ to $^{36}\text{S}$ (tensor)	From $^{36}\text{S}$ to $^{34}\text{Si}$ (spin orbit)
Splitting	<b>D1S</b>	<b>D1S</b>
$2p$	13%	43%
	<b>D1ST2a</b>	<b>D1ST2a</b>
$2p$	40%	39%
	<b>D1ST2c</b>	<b>D1ST2c</b>
$2p$	27%	42%

Reductions of the neutron  $2p$  splitting.

M. Grasso and M. Anguiano, Phys. Rev. **C92**, 054216 (2015)

# Interplay between tensor force and pairing correlations

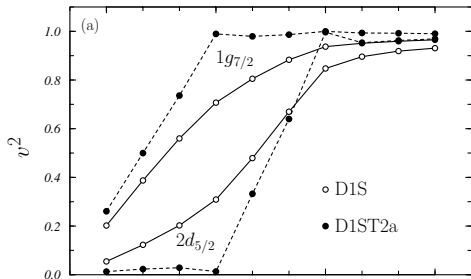


Pairing field in HFB:  $V_{\text{Central}} + V_{\text{SO}} + V_{\text{Coul}}$ .

Pairing field in BCS:  $V_{\text{Central}}$

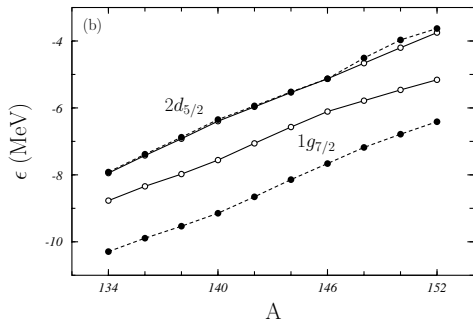
# Interplay between tensor force and pairing correlations

$N = 82$



Proton Levels

Tensor force?



neutron  $1h_{11/2}$



# A new fit for the tensor interaction

$$V_{\text{tensor}}(\vec{r}_1, \vec{r}_2) = V_T S_{12} e^{-(r_1-r_2)^2/\mu_T^2} + V_{T\tau} \tau(1) \cdot \tau(2) S_{12} e^{-(r_1-r_2)^2/\mu_{T\tau}^2}$$

1. A more general fit, with five parameters

$$W_{\text{LS}} \quad V_T \quad V_{T\tau} \quad \mu_T \quad \mu_{T\tau}$$

2. How? Using the procedure following to obtain D1ST2c, and adding two observables more.

**Energy of the first  $0^-$  state of  $^{16}\text{O}$  and  $1^+$  GT of  $^{48}\text{Ca}$**